

# Non-homogeneous incompressible Bingham flows with variable yield stress and application to volcanology.

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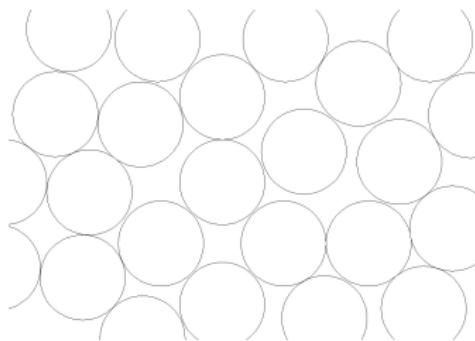
june 2015





# Laboratory experiment

Zoom into the granular column →

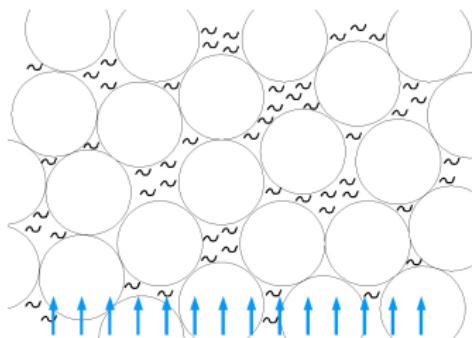


O. ROCHE, LMV.

# Laboratory experiment

Zoom into the granular column →

Fluidisation :  
injection of **gas**  
through  
a pore plate.



O. ROCHE, LMV.

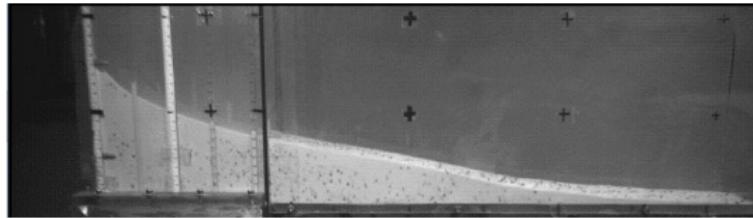


Figure: Non-fluidised → **short** runout distance



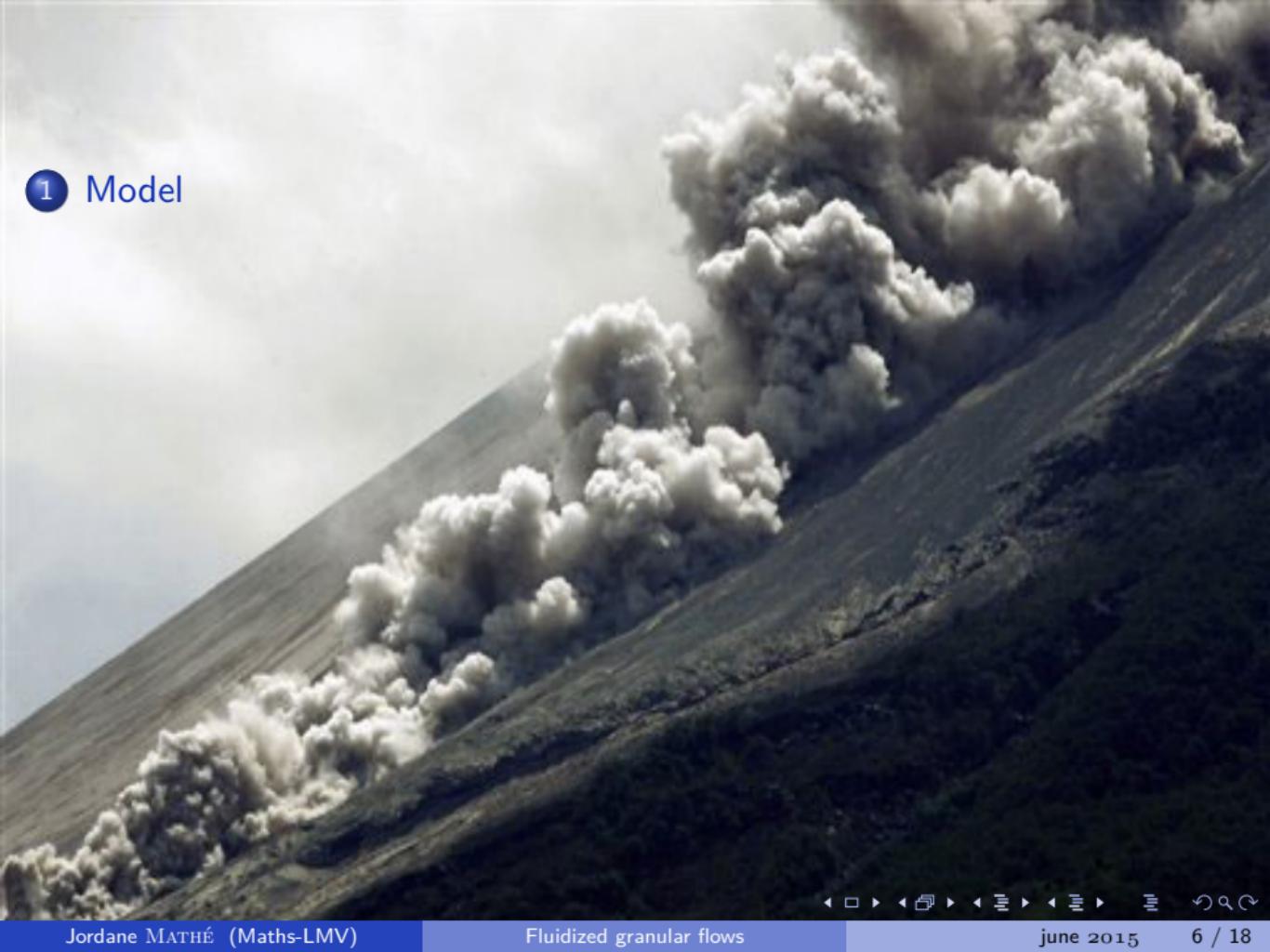
Figure: Fluidised → **long** runout distance

## 1 Model

## 2 Numerical simulation

## 3 Perspectives

# 1 Model



## Two phases for one fluid

Consider **one mixed fluid** with **variable density**.

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$$\begin{cases} \operatorname{div}(v) = 0 \\ \partial_t \rho + \operatorname{div}(\rho v) = 0 \\ \partial_t(\rho v) + \operatorname{div}(\rho v \otimes v) + \nabla p = \rho g + \operatorname{div}(\mathbf{S}) \end{cases}$$

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Unknown:

- $v$ : velocity,
- $p$ : total pressure,
- $\rho$ : density.

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## Rheology

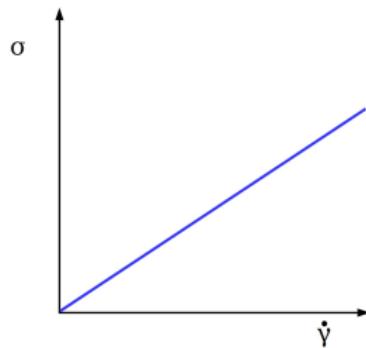
Let precise  $\mathbf{S}$ .

# Rheology

In our case

$$\mathbf{S} = \mu D\mathbf{v} +$$

where  $D\mathbf{v} = \dot{\gamma}$  is the strain rate tensor,  
 $\mu$  is the effective viscosity



# Rheology

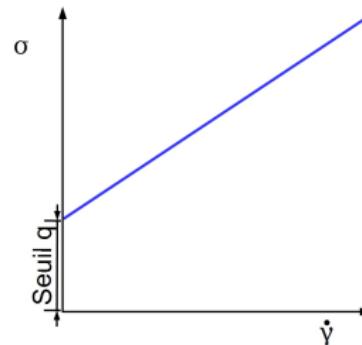
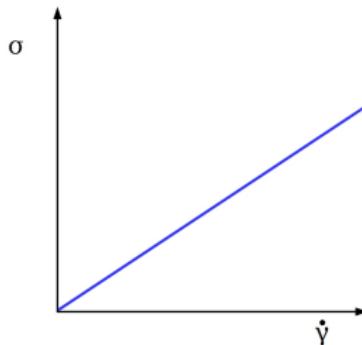
In our case

$$\mathbf{S} = \mu Dv + \Sigma,$$

where  $Dv = \dot{\gamma}$  is the strain rate tensor,  
 $\mu$  is the effective viscosity

$$\Sigma = q \frac{Dv}{|Dv|}$$

$\Rightarrow$  Bingham fluid with yield stress  $q$



# Rheology

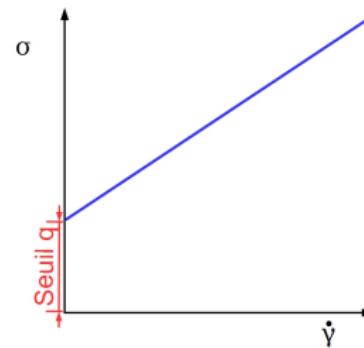
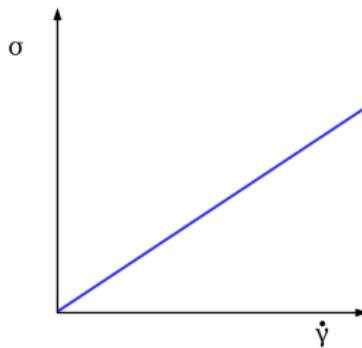
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## Idea

Let vary the yield stress  $q$  as a function of the **interstitial gas pressure**.

# Variation of the yield stress

## Definition of the yield stress

$$q = \begin{cases} \text{atmospheric pressure:} & \text{Coulomb friction} \\ \text{high pressure:} & \text{fluid} \end{cases}$$

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where  $\delta$  is the internal friction angle.

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$$q = \tan(\delta)(\rho gh - \text{gas pressure})^+$$

where  $\delta$  is the internal friction angle.

## Equations of the model

Noting  $p_f$  the interstitial gas pressure, we obtain:

$$\left\{ \begin{array}{l} \operatorname{div}(v) = 0 \\ \partial_t \rho + \operatorname{div}(\rho v) = 0 \\ \partial_t(\rho v) + \operatorname{div}(\rho v \otimes v) - \mu \Delta v + \nabla p = \tan(\delta) \operatorname{div} \left( \underbrace{(\rho g h - p_f)^+}_{\text{yield } q} \frac{Dv}{|Dv|} \right) + \rho g \end{array} \right.$$

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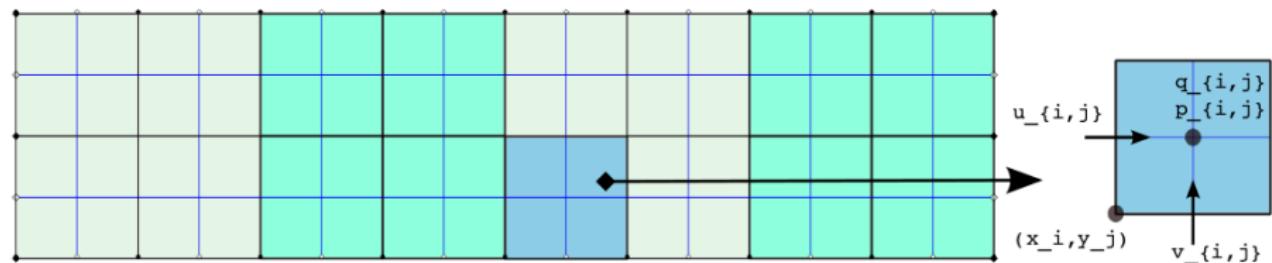
where  $\kappa$  is the diffusion coefficient.

A photograph of a volcano erupting, with a massive column of white smoke and ash rising into a clear blue sky. The base of the volcano is visible, showing green vegetation and some rocky terrain.

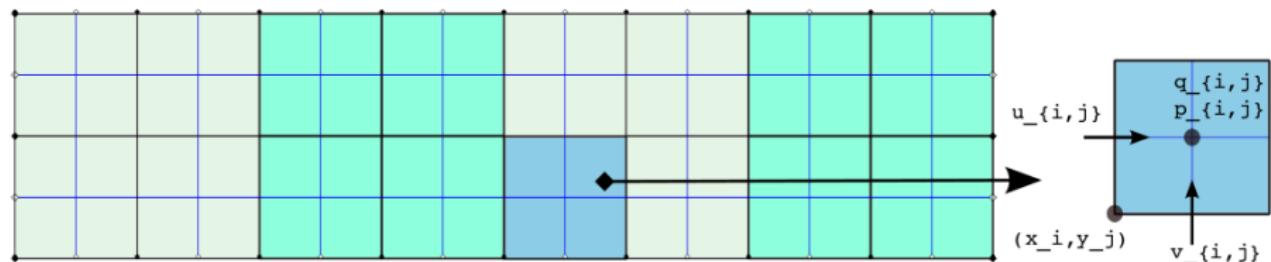
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## Numerical simulation

# Dambreak: Mesh



# Dambreak: how to treat the density?



$$\partial_t \rho + \mathbf{v} \cdot \nabla \rho = 0$$

numerical method:

RK3 TVD - WENO5 scheme.

## Numerical scheme (without density)

$n = 0$        $v^0$ ,  $\Sigma^0$ ,  $p^0$  and  $p_f^0$  given.

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We compute  $(v^{n+1}, p^{n+1})$  thanks to the incompressibility constrain

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# Experimental conditions

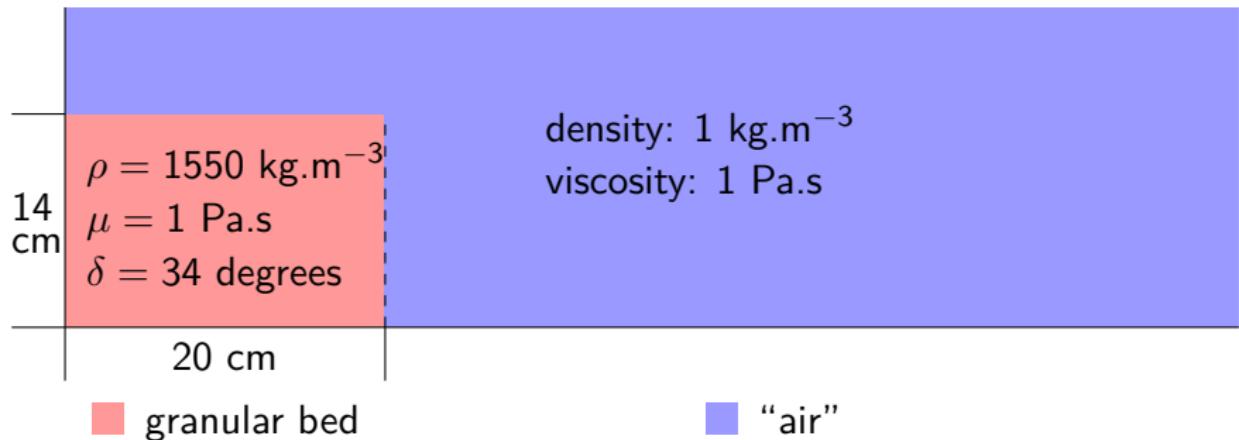


Figure: Experimental setup

# Numerical simulation

We compute the collapse with different diffusion coefficient  $\kappa$ .

- With  $\kappa = 1$ , it happens nothing.
- $\kappa = 0.2$
- $\kappa = 0.1$

Diffusion coefficient = 0.2

Diffusion coefficient = 0.1



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## Perspectives

# Perspectives

- To compare the results given by this numerical scheme for the **non-fluidised** case:
  - ▶ Lower yield stress, . . .
  - ▶ Lower friction coefficient . . .
- To compare the results given by this numerical scheme for the **fluidised** case:
  - ▶ Adapt the viscosity value.

# Thank you!

