

Stein's method for Gibbs point processes

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joint work with Dominic Schuhmacher

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A short introduction to Stein's method

Theorem (Stein's Lemma)

Let $Z \sim \mathcal{N}(0, 1)$. Then

$$\mathbf{E}f'(Z) - \mathbf{E}Zf(Z) = 0$$

for all functions such that the above expectations exist. Conversely, every random variable satisfying this equation for a large enough class of functions f is necessarily the standard normal distribution.

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- The operator $(\mathcal{A}f)(x) = f'(x) - xf(x)$ characterises the normal distribution.

- The *Stein equation*

$$f(x) - \mathbf{E}f(Z) = \mathcal{A}h_f(x)$$

is solved by the *Stein solution*

$$h_f(x) = e^{\frac{x^2}{2}} \int_{-\infty}^x (f(y) - \mathbf{E}f(Z)) e^{-\frac{y^2}{2}} dy.$$

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- Assume we want to bound $\sup_{f \in \mathcal{F}} |\mathbf{E}f(X) - \mathbf{E}f(Z)|$. Using the Stein equation we can bound $\sup_{f \in \mathcal{F}} |\mathbf{E}\mathcal{A}h_f(X)|$ instead.
- The *Stein factors* $\|h'_f\|$ and $\|h''_f\|$ play a crucial role in bounding $\mathbf{E}\mathcal{A}h_f(X)$.

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- Every $\xi \in \mathfrak{N}$ can be written as $\xi = \sum_{i=1}^n \delta_{x_i}$ for some $x_1, \dots, x_n \in \mathcal{X}$, and where δ_x denotes the Dirac measure at point x .

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In particular $\mathbb{E}(\Xi(A)) = \lambda(A)$.
 - (b) *The numbers of points in any two disjoint sets are independent:*
 $\Xi(A), \Xi(B)$ independent for any $A, B \in \mathcal{B}$ with $A \cap B = \emptyset$.
- We write $\Xi \sim \text{PoP}(\lambda)$.

- A function $u: \mathfrak{N} \rightarrow \mathbb{R}_+$ is called **hereditary** if $u(\xi) = 0$ implies $u(\eta) = 0$ for all point configurations $\xi, \eta \in \mathfrak{N}$ with $\xi \subset \eta$.

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- A Gibbs process is completely described by its **conditional intensity** $\lambda(\cdot | \cdot)$, where

$$\lambda(x | \xi) = \frac{u(\xi + \delta_x)}{u(\xi)} \quad \text{for all } \xi \in \mathfrak{N}, x \in \mathcal{X} \text{ with } \xi(\{x\}) = 0.$$

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- We write $\text{Gibbs}(\lambda)$ for the distribution of this Gibbs process.

Examples of Gibbs processes I

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- A Gibbs process is a *pairwise interaction process* (PIP) if the conditional intensity is of the form $\lambda(x | \xi) = \beta(x) \prod_{y \in \xi} \varphi(x, y)$, for a $\beta: \mathcal{X} \rightarrow \mathbb{R}_+$ and a symmetric *interaction function* $\varphi: \mathcal{X} \times \mathcal{X} \rightarrow \mathbb{R}_+$, e.g. for the *Strauss process*

$$\varphi(x, y) = \begin{cases} \gamma & \text{if } d_0(x, y) \leq r; \\ 1 & \text{otherwise,} \end{cases}$$

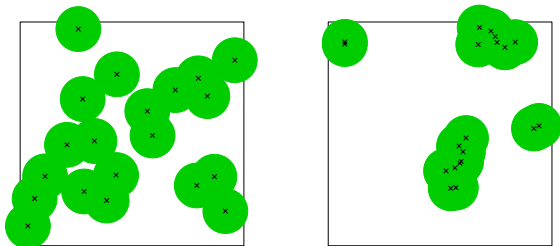
for a $r > 0$ and a $\gamma \in [0, 1]$.

Examples of Gibbs processes II

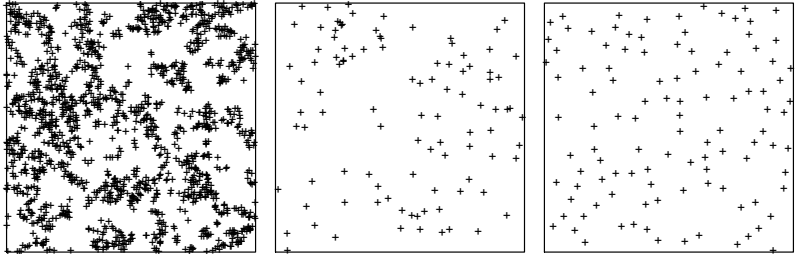
- The *area-interaction process (AIP)* has the conditional intensity

$$\lambda(x \mid \xi) = \beta \gamma^{-\alpha(U_r(\xi + \delta_x) \setminus U_r(\xi))},$$

where $\beta, \gamma, r > 0$ and $U_r(\xi) = \bigcup_{x \in \xi} \mathbb{B}_r(x)$ denotes the green area.



Simulated Gibbs processes



Left: Aip with $\gamma = 100$, Middle: Aip with $\gamma = 0.01$, Right: Strauss process with $\gamma = 0$

Spatial birth-death processes

- Suppose that we have birth rates and death rates

$$b(\cdot | \cdot): \mathcal{X} \times \mathfrak{N} \rightarrow \mathbb{R}_+ \quad \text{with } \bar{b}(\xi) := \int b(x | \xi) \alpha(dx) < \infty;$$

$$d(\cdot | \cdot): \mathcal{X} \times \mathfrak{N} \rightarrow \mathbb{R}_+ \quad \text{with } \bar{d}(\xi) := \sum_{x \in \xi} d(x | \xi) < \infty.$$

- Let $\bar{a}(\xi) = \bar{b}(\xi) + \bar{d}(\xi)$.

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- Let $\bar{a}(\xi) = \bar{b}(\xi) + \bar{d}(\xi)$.
- A $\text{SBD}^{(\xi_0)}(b, d)$ -process is a pure-jump Markov process on \mathfrak{N} that starts in $\xi_0 \in \mathfrak{N}$ and holds each state ξ for an $\text{Exp}(\bar{a}(\xi))$ -distributed time, after which
 - (a) with probability $\bar{b}(\xi)/\bar{a}(\xi)$ a point is added, positioned according to the density $b(\cdot | \xi)/\bar{b}(\xi)$, or
 - (b) with probability $d(x | \xi)/\bar{a}(\xi)$ the point at x is deleted.

- In what follows always $b(\cdot | \cdot) = \lambda(\cdot | \cdot)$, $d \equiv 1$ (“unit per-capita death rate”) and $\lambda(\cdot | \cdot)$ is *locally stable*, i.e. there exists a constant c^* such that

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- Let $Z = (Z(t))_{t \geq 0} \sim \text{SBD}(\lambda, 1)$. Then
- Z is non-explosive;
- Z has $\text{Gibbs}(\lambda)$ as its unique stationary distribution.
- Z has the infinitesimal generator

$$\mathcal{A}h(\xi) = \int_{\mathcal{X}} [h(\xi + \delta_x) - h(\xi)] \lambda(x | \xi) \alpha(dx) + \int_{\mathcal{X}} [h(\xi - \delta_x) - h(\xi)] \xi(dx)$$

for certain functions $h: \mathfrak{N} \rightarrow \mathbb{R}$.

- Our goal is to define a coupling of two Z and \tilde{Z} SBD($\lambda, 1$)'s which follow the same dynamics but are started in different configurations.
- Given that at time t the processes are in states $Z(t) = \xi$ and $\tilde{Z} = \tilde{\xi}$, propose a birth with rate $\max(\lambda(\cdot | \xi), \lambda(\cdot | \tilde{\xi}))$. The first process Z accepts it with probability $\lambda(\cdot | \xi) / \max(\lambda(\cdot | \xi), \lambda(\cdot | \tilde{\xi}))$ and \tilde{Z} accepts it with probability $\lambda(\cdot | \xi) / \max(\lambda(\cdot | \tilde{\xi}), \lambda(\cdot | \tilde{\xi}))$.

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- The deaths are coupled in the obvious manner, each point dies with rate one independently of the others but of course the common points of Z and \tilde{Z} die together.

- Let

$$\varepsilon = \sup_{\|\xi - \eta\|=1} \int_{\mathcal{X}} |\lambda(x | \xi) - \lambda(x | \eta)| \alpha(dx) \quad \text{and} \quad c = c^* \alpha(\mathcal{X}).$$

- Consider the *coupling time* $\tau = \inf\{t \geq 0: Z_t^\xi = \tilde{Z}_t^{\xi + \delta_x}\}$.

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Theorem

We have $\mathbf{E}\tau < \infty$. In particular if $\varepsilon < 1$ then $\mathbf{E}\tau < (1 + \varepsilon)/(1 - \varepsilon)$.

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Example

For a Strauss processes in \mathbb{R}^d with parameters β, γ, R we get $\varepsilon = \beta(1 - \gamma)\alpha_d R^d$, where α_d is the volume of the unit ball in \mathbb{R}^d .

Our goal:

Find upper bound for the total variation distance

$$d_{TV}(\text{Gibbs}(\nu), \text{Gibbs}(\lambda)),$$

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A brief history:

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- **Barbour and Brown (1992)**. Poisson process approximation.

Set-up for general probability metrics

- Let $H \sim \text{Gibbs}(\lambda)$. Suppose we want to bound

$$d(\mathcal{L}(\Xi), \text{Gibbs}(\lambda)) = \sup_{f \in \mathcal{F}} |\mathbb{E}f(\Xi) - \mathbb{E}f(H)|,$$

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- For the total variation metric

$$\mathcal{F} = \mathcal{F}_{TV} = \{1_C; C \in \mathcal{N}\}.$$

Setting up the Stein equation

- For every $f \in \mathcal{F}$ find $h = h_f: \mathfrak{N} \rightarrow \mathbb{R}$ such that

$$f(\xi) - \mathbb{E}f(\mathbf{H}) = \mathcal{A}h_f(\xi) \quad \text{for all } \xi \in \mathfrak{N}, \quad (\text{Stein equation})$$

where \mathcal{A} is the generator of a Markov process with stationary distribution $\text{Gibbs}(\lambda)$ (generator approach).

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- Natural choice: the $\text{SBD}^{(\xi)}(\lambda, 1)$ -process $Z^{(\xi)} := (Z_t^{(\xi)})_{t \geq 0}$ from earlier.

Solution of the Stein equation

- It can be shown that for bounded f the function $h = h_f : \mathfrak{N} \rightarrow \mathbb{R}$,

$$h(\xi) := - \int_0^\infty [\mathbb{E}f(Z_t^{(\xi)}) - \mathbb{E}f(H)] dt,$$

is well-defined and solves the Stein equation

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- To bound $\sup_{f \in \mathcal{F}_{TV}} |\mathbf{E} \mathcal{A} h_f(\Xi)|$ it turns out that the key ingredient is the *Stein factor*

$$c_1(\lambda) = \sup_{f \in \mathcal{F}_{TV}, x \in \mathcal{X}, \xi \in \mathcal{N}} |h_f(\xi + \delta_x) - h_f(\xi)|.$$

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- Note that

$$h_f(\xi + \delta_x) - h_f(\xi) = - \int_0^\infty [\mathbb{E} f(Z_t^{(\xi)}) - \mathbb{E} f(Z_t^{(\xi + \delta_x)})] dt.$$

- Thus $c_1(\lambda)$ can be bounded by the expected coupling time of two SBDP's starting in configurations differing by one point.

Theorem (Schuhmacher and S, 2012)

For any two Gibbs point processes

Ξ with conditional intensity $\nu(\cdot | \cdot)$,

H with locally stable conditional intensity $\lambda(\cdot | \cdot)$,

there exists a finite constant $c_1(\lambda)$ such that

$$d_{TV}(\mathcal{L}(\Xi), \mathcal{L}(H)) \leq c_1(\lambda) \int_{\mathcal{X}} \mathbb{E} |\nu(x | \Xi) - \lambda(x | \Xi)| \alpha(dx).$$

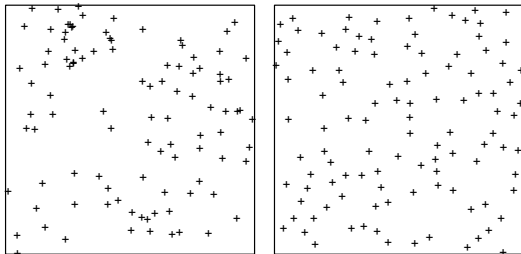
The Stein factor $c_1(\lambda)$ is bounded by the expected coupling time from earlier.

Pairwise interaction processes

- Suppose that $\mathcal{X} \subset \mathbb{R}^d$, and $\Xi \sim \text{PIP}(\beta, \varphi_1)$ and $H \sim \text{PIP}(\beta, \varphi_2)$ are stationary and inhibitory, i.e. β is constant and $\varphi_i(x, y) = \varphi_i(x - y) \leq 1$ for all $x, y \in \mathcal{X}$. Then

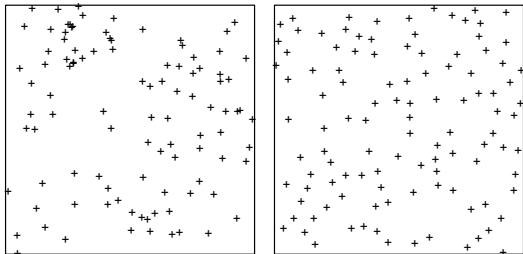
$$d_{\text{TV}}(\mathcal{L}(\Xi), \mathcal{L}(H)) \leq c_1(\lambda) \beta \mathbb{E}|\Xi| \int_{\mathbb{R}^d} |\varphi_1(x) - \varphi_2(x)| dx.$$

Convergence of the Area interaction process



Left: Aip with $\gamma = 0.01$, Right: Strauss process with $\gamma = 0$

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- Suppose that $\mathcal{X} \subset \mathbb{R}^d$, and $\Xi \sim \text{AIP}(\beta\gamma^{\alpha_d(R/2)^d}, \gamma; R/2)$ and $H \sim \text{Strauss}(\beta, 0; R)$, where α_d is the volume of the unit ball in \mathbb{R}^d . Then

$$d_{\text{TV}}(\mathcal{L}(\Xi), \mathcal{L}(H)) \leq c_1(\lambda) 2d\alpha_d R^{d-1} \beta \mathbb{E}|\Xi| (\log \gamma^{-\alpha_d})^{-1/d}.$$

Non locally stable Gibbs processes

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Non locally stable Gibbs processes

- Although most of the Gibbs processes considered in spatial statistics are locally stable, there exist some notable exceptions (e.g. the *Lennard - Jones process*).
- For pairwise interaction processes satisfying some mild assumptions, the previous Theorem can be generalised to

$$d_{\text{TV}}(\mathcal{L}(\Xi), \mathcal{L}(\text{H})) \leq C_1 \|\varphi_1 - \varphi_2\|_{L^1} + C_2(C_1),$$

where the constant C_2 can be chosen arbitrarily small causing a larger C_1 .

The probability generating functional

- Assume that $\mathcal{F} = \{f\}$ consists of only one function, namely $f(\xi) = \prod_{x \in \xi} g(x)$. Then

$$\Psi_{\Xi}(g) = \mathbf{E}\left(\prod_{x \in \Xi} g(x)\right)$$

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is called probability generating functional.

- Assume that our Gibbs processes live on a subset of \mathbb{R}^d and that the H is a homogeneous Poisson process with intensity ν . Then from the Stein equation we get

$$\Psi_{\Xi}(g) - \exp\left(-\nu \int_{\mathbb{R}^d} 1 - g(x) dx\right) = \mathbf{E}(\mathcal{A}h_f(\Xi)).$$

Theorem

Let Ξ be a stationary and locally stable Gibbs point process on \mathbb{R}^d with intensity λ and local stability constant c^* . Then

$$1 - \lambda G \leq \Psi_{\Xi}(g) \leq 1 - \frac{\lambda}{c^*} (1 - e^{-c^* G}),$$

where $G = \int_{\mathbb{R}^d} 1 - g(x) dx$.

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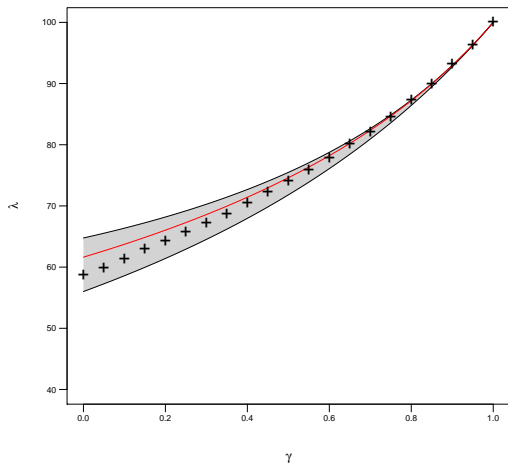
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Theorem

Let $\Xi \sim \text{PIP}(\beta, \varphi)$ be inhibitory, with finite interaction range, i.e. $1 - \varphi$ has bounded support, and with intensity λ . Then

$$\frac{\beta}{1 + \beta G} \leq \lambda \leq \frac{\beta}{2 - e^{-\beta G}},$$

where $G = \int_{\mathbb{R}^d} 1 - \varphi(x) dx$.



The graphic shows the intensities of Strauss processes in \mathbb{R}^2 with $\beta = 100$ and $r = 0.05$. The crosses are simulated values, the red line is the PS-approximation and the grey area corresponds to our bounds.



D. Schuhmacher and K. Stucki.

Gibbs point process approximation: Total variation bounds using stein's method.

Annals of Probability, 42(5):1911–1951, 2014



K. Stucki and D. Schuhmacher.

Bounds for the probability generating functional of a Gibbs point process.

Adv. in Appl. Probab., 46:21–34, 2014